



## Quantencomputing und Quantensimulation

### Wintersemester 2023 - Übungsblatt 10

Ausgabe: 19.01.2024, Abgabe: 26.01.2024, Übungen: 29.01.2024

#### Aufgabe 24: Jordan-Wigner Transformation (8 points)

In the lecture, the Jordan-Wigner transformation was introduced to transform fermionic systems into spin systems, which can be simulated on a quantum computer. The vacuum state (the state without particles) is simulated as a spin up state and the one-particle state as a spin down state,

$$|0\rangle = a|1\rangle \equiv |\uparrow\rangle, \quad |1\rangle = a^\dagger|0\rangle \equiv |\downarrow\rangle.$$

At first sight, this gives us an equivalence between the ladder operators of the spin states  $\sigma_\pm = 1/2(\sigma_x \pm i\sigma_y)$  and the creation and annihilation operators of the fermions  $a^{(\dagger)}$ ,

$$a \equiv \sigma_+, \quad a^\dagger \equiv \sigma_-.$$

a) (1 point) Use  $[\sigma_+, \sigma_-] = \sigma_z$  to show  $1 - 2a^\dagger a \equiv \sigma_z$ .

Let us now consider fermions on a chain. This shows that the above equivalence no longer works, as the creation and annihilation operators for different lattice sites ( $i \neq j$ ) has to anti-commute, however for different lattice sites the ladder operators commute,

$$a_i^\dagger a_j^\dagger = -a_j^\dagger a_i^\dagger \quad \leftrightarrow \quad \sigma_{i,-} \sigma_{j,-} = \sigma_{j,-} \sigma_{i,-}.$$

In order to express this property using the spin operators, the transformation

$$a_i \rightarrow \sigma_{i,+} \otimes_{k=1}^{i-1} \sigma_{k,z}, \quad a_i^\dagger \rightarrow \sigma_{i,-} \otimes_{k=1}^{i-1} \sigma_{k,z}$$

is performed. We therefore accumulate a phase of  $-1$  per spin down state on the lattice sites before  $i$ .

b) (2 points) Show that this preserves the canonical anti-commutation relations  $\{a_i, a_j^\dagger\} = \delta_{i,j}$  and  $\{a_i, a_j\} = 0$ .

c) (1 point) Show that  $a_n^\dagger a_{n+1} \equiv \sigma_{n,-} \sigma_{n+1,+}$ .

d) (1 point) Use the Jordan-Wigner transformation on the following fermionic Hamiltonian

$$H = \sum_n J_z \left( 1 - 2(a_n^\dagger a_n + a_{n+1}^\dagger a_{n+1}) - 4a_n^\dagger a_{n+1}^\dagger a_n a_{n+1} \right) + \frac{J_\perp}{2} \left( a_n^\dagger a_{n+1} + a_{n+1}^\dagger a_n \right)$$

to obtain

$$H = \sum_n J_z \sigma_{n,z} \sigma_{n+1,z} + \frac{J_\perp}{2} (\sigma_{n,-} \sigma_{n+1,+} + \sigma_{n,+} \sigma_{n+1,-}).$$

The Jordan-Wigner transformation uses the occupation number representation. An alternative formulation is given by using the so-called parity basis. In the parity basis, the state of the  $i$ th qubit  $q_i$  is given by the formula  $q_i = \sum_{j < i} f_j \bmod 2$ , where  $f_i$  describes the state of the  $i$ th qubit in the occupation number representation ( $f_i = 1$  if a fermion occupies lattice site  $i$ ). This alternative transformation is obtained as

$$a_i \rightarrow \bigotimes_{k=0}^{M-i-1} \sigma_{M-k,x} \sigma_{i,x} \sigma_{i-1,z} + i \bigotimes_{k=0}^{M-i-1} \sigma_{M-k,x} \sigma_{i,y},$$

$$a_i^\dagger \rightarrow \bigotimes_{k=0}^{M-i-1} \sigma_{M-k,x} \sigma_{i,x} \sigma_{i-1,z} - i \bigotimes_{k=0}^{M-i-1} \sigma_{M-k,x} \sigma_{i,y}.$$

d) (1 point) Write the state  $|10100111\rangle$  given in the occupation number representation in the parity basis.

e) (2 points) Show that this transformation also preserves the canonical anti-commutation relations  $\{a_i, a_j^\dagger\} = \delta_{i,j}$  and  $\{a_i, a_j\} = 0$ .

### Aufgabe 25: Simulation of a one dimensional particle (3 points)

We consider the simulation of a one dimensional particle with the Hamiltonian

$$H = \frac{1}{2}p^2 + V(x).$$

In order to simulate this Hamiltonian with a circuit, we discretize space by using the (dimensionless) position operator  $x = \sum_x x |x\rangle\langle x|$  and the approximated (dimensionless) momentum operator  $p = -\frac{i}{2} \sum_x (|x+1\rangle\langle x| - |x-1\rangle\langle x|)$  ( $x$  here represents a binary number and  $x\Delta x = xL/2^n$  describes the actual location). A general state  $|\psi\rangle$  can then be described by  $|\psi\rangle = \sum_x \psi(x\Delta x) |x\rangle$ , where  $\psi(x\Delta x)$  represents the wave function.

a) (1 point) Given the state  $|p\rangle = U_{\text{QFT}} |x\rangle$  with  $U_{\text{QFT}} = \frac{1}{\sqrt{2^n}} \sum_y e^{2\pi i xy/2^n} |y\rangle\langle x|$ . Calculate  $p|p\rangle$  for both the exact operator  $p = -i\partial_x$  and its discretized approximation and compare the results.

b) (2 points) To implement the effect of the potential, the unitary mapping

$$U_V : |x\rangle |y\rangle \rightarrow |x\rangle |y \oplus \Delta t V(x)\rangle$$

is used. Show that the circuit shown below produces the state

$$U_V |\psi(0)\rangle U_{\text{QFT}}^\dagger |1\rangle = \sum_{x=1}^{2^n-1} \langle x | \psi(0) \rangle e^{-2\pi i \Delta t V(x)/2^t} |x\rangle U_{\text{QFT}}^\dagger |1\rangle.$$

